Influence of radiation defects on exciton-magnon interactions in nickel oxide

N. Mironova-Ulmane^{*a}, V. Skvortsova^a, A. Kuzmin^a, I. Sildos^b ^a Institute of Solid State Physics, University of Latvia, LV-1063 Riga, Latvia ^b Institute of Physics, Riia street 142, EE-2400 Tartu, Estonia

ABSTRACT

Influence of radiation defects on the optical absorption spectrum of nickel oxide (NiO) was studied at the temperature 6 K in the near infra-red energy range 7750-8300 cm⁻¹, corresponding to the magnetic-dipole transition ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ at nickel sites. NiO single-crystals, grown by the method of chemical transport reactions on MgO(100) substrates, were irradiated by the neutron fluences up to 5×10^{18} cm⁻². Two sharp lines were observed at the low-energy side of the band: the peak at 7805 cm⁻¹ is assigned to the pure exciton transition, whereas the peak at 7845 cm⁻¹ to the exciton-magnon excitation, which occurs at the Brillouin zone-centre (BZC). An increase of the defects concentration at higher fluences results in a lowering of the magnon satellite peak intensity. The long-wavelength BZC magnon absorption is sensitive to the long-range magnetic ordering, which becomes destroyed in the presence of radiation defects. Therefore, the observed decrease of the peak intensity is attributed to a decrease of the spin-spin correlation length due to inhomogenious broadening.

Keywords: NiO, magnetic-dipole transition, radiation defects, magnons

1. INTRODUCTION

The electronic and magnetic structure of nickel oxide (NiO) is a subject of continuous interest during many years [1]. NiO is a type-II easy-plane antiferromagnet with the Néel temperature $T_N = 523$ K [2]. In the paramagnetic phase, it has a rocksalt-type crystal structure, whereas a week cubic-to-rhombohedral distortion occurs below T_N in the antiferromagnetic phase [3]. The antiferromagnetic structure of NiO is characterized by dominating superexchange interactions ($J_{NNN} \approx 150$ cm⁻¹) between the next-nearest-neighbours (NNN), connected by 180° Ni²⁺-O²⁻-Ni²⁺ linear atom chains. Much weaker superexchange coupling ($|J_{NN}| \approx 11$ cm⁻¹) occurs between the nearest-neighbours (NN), connected by 90° Ni²⁺-O²⁻-Ni²⁺ paths [4-6]. The exchange interactions in NiO not only influence its magnetic properties but also appear in optical absorption [7-11], photoluminescence [12], and Raman scattering [4,6,13-17] spectra.

The magnetic properties of nickel oxide can be strongly modified by the presence of impurity atoms and radiation defects [18]. For example, substitution of a part of nickel ions by magnesium ions results in a continuous series of Ni_cMg_{1-c}O solid solutions, having a rich magnetic phase diagram [3,19]. At T = 0 K, it consists of four regions: (1) a homogeneous antiferromagnet ($0.63 \le c \le 1$), (2) a cluster antiferromagnet ($0.4 \le c < 0.63$), (3) spin-glass type ($0.25 \le c < 0.4$), and (4) a paramagnet ($c \le 0.2$). A number of other mixed systems as Li_xNi_{1-x}O [20], Co_xNi_{1-x}O [21,22] and Fe_xNi_{1-x}O [23] were also studied. The influence of neutron irradiation on optical absorption spectra in NiO and Ni_cMg_{1-c}O solid solutions was studied in [18]: it was found that irradiation increases mainly the intensity of spin-forbidden optical transitions.

In the present study we will concentrate on the optical absorption. The absorption bands observed in nickel oxide are related to parity-forbidden d-d transitions and can be labelled using the energy level diagram of a free nickel ion Ni²⁺(3d⁸) in a cubic crystal field. Three transitions from the ${}^{3}A_{2g}(F)$ ground state to the ${}^{3}T_{2g}(F)$, ${}^{3}T_{1g}(F)$ and ${}^{3}T_{1g}(P)$

^{*} <u>ulman@latnet.lv</u>

excited states are spin-allowed ($\Delta S=0$), but the others are spin-forbidden. The bands in UV and visible spectrum ranges are the electric-dipole in nature, whereas the near-IR band at 8800 cm⁻¹ corresponds to the magnetic-dipole transition ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ [18]. We have found previously [7,9-11] that the last band contains at low temperatures two zero-phonon lines, located at about 7805 cm⁻¹ and 7845 cm⁻¹, due to the pure exciton and exciton–one-magnon excitations, respectively. The intensity of the exciton–one-magnon line varies strongly both with temperature [10,11] and composition [9,11], due to homogeneous and inhomogeneous broadening, respectively.

In this work, we discuss for the first time an influence of the radiation defects, produced by neutrons, on the exciton– one-magnon zero-phonon transition within the magnetic-dipole ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ band in NiO.

2. EXPERIMENTAL

Single-crystal NiO(100) were grown epitaxially on freshly cutted and polished single-crystal MgO(100) substrates by chemical transport reactions method (the "sandwich" technique) using HCl gas [24]. Thus obtained NiO crystals have green colour and retain orientation of the MgO substrate. Several samples of NiO were irradiated at the Salaspils (Latvia) nuclear reactor by the fast neutrons with the energy >0.1 MeV and fluences 1.8×10^{18} and 5×10^{18} cm⁻². The cadmium filter was used to eliminate thermal neutrons. The irradiation was performed by placing the samples into the water channel, so that the samples temperature during irradiation was maintained below 400 K. Optical absorption spectra were recorded using the split-beam Jasco spectrophotometer (Model V-570) with the tungsten iodine lamp used as a source. PbS photoconductive cell was used as a detector. A liquid helium cryostat was used to control the temperature of the samples down to 5 K with an accuracy ±1 K.



Figure 1: Temperature dependence of the ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ absorption band in NiO(100) single-crystal. Peak A corresponds to pure exciton absorption, whereas peak B to the exciton–one-magnon excitation. The fine structure above 7900 cm⁻¹ is due to the exciton–phonon excitations.

3. RESULTS AND DISCUSSION

The magnetic-dipole band ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ in NiO consists at T=10 K of two sharp zero-phonon lines (peaks A and B in Fig. 1), phonon satellite peaks and a broad sideband [7,10]. The peak A at ~7805 cm⁻¹ is attributed to the pure exciton transition, whereas the peak B at ~7845 cm⁻¹ to the simultaneous exciton–one-magnon excitation [7,10]. Here the one-magnon excitation occurs at the Brillouin zone-centre, and its energy can be obtained as the difference between peaks A and B [10]. Upon increasing temperature or creation of defects under neutron fluence, the ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ band experiences homogeneous and inhomogeneous broadening, which affects the intensity and the position of the two zero-phonon peaks and of the phonon sideband [25].

Temperature dependence of the one-magnon sideband in NiO is shown in Fig. 1. Note that the peak B disappears above 110 K that is well below the Néel temperature $T_N = 523$ K. The separation between two peaks A and B remains constant at all measured temperatures and equals to about 41 cm⁻¹. This value is in agreement with that obtained previously by far-infrared antiferromagnetic resonance (AFMR) [26,27] and by Raman scattering [14-17]. A small shift of both (A and B) peaks as well as of the band maximum to smaller energies, observed upon temperature increase, is attributed to the change of the crystal field strength due to the thermal lattice expansion [18].



Figure 2: The ${}^{3}A_{2g}(F) \rightarrow {}^{3}T_{2g}(F)$ absorption band in NiO(100) singlecrystals, measured at the temperature T = 6 K. The two lower curves correspond to samples, irradiated by the neutron fluences 1.8×10^{18} and 5×10^{18} cm⁻². They contain a sharp peak at ~9550 cm⁻¹, associated with the radiation defects Fe²⁺ in MgO substrate. Peak A corresponds to pure exciton absorption, whereas peak B to the exciton–one-magnon excitation. The fine structure above 7900 cm⁻¹ is due to the exciton–phonon excitations.

Previous investigations [18] of the influence of neutron irradiation on optical absorption spectra of pure nickel oxide and $Ni_cMg_{1-c}O$ solid solutions have shown that the irradiation increases intensity of spin-forbidden transitions and does not affect the intensity of spin-allowed transitions. An increase of the spin-forbidden transitions intensity occurs due to the exchange interaction between 3d electrons of nickel ion and the electron, localized at a defect site [18,28]. Such interaction should produce opposite effect, i.e. results in a decrease of intensity for the magnon assisted absorption, since the presence of radiation defects will destroy periodic magnetic ordering. Previously we have found that a substitutional modification of the magnetic sublattice by an admixture of less than 10% of magnesium ions to nickel oxide leads to a strong damping of the peak B intensity already at T = 6 K [3,19]. As one can see in Fig. 2, irradiation of nickel oxide with neutrons leads to a lowering of the peak B intensity, which decreases with an increase of the neutron fluence due to inhomogeneous broadening produced by the presence of radiation defects.

Besides, a sharp peak at 9550 cm⁻¹ and a fine structure above it are observed for our samples after irradiation. Their intensity increases with the neutron fluence variation from 1.8×10^{18} cm⁻² to 5×10^{18} cm⁻². We associate the peak at 9550 cm⁻¹ and the fine structure above it with the transition ${}^{5}T_{2g} \rightarrow {}^{5}E_{g}$ in the radiation induced defects Fe²⁺ present in MgO substrate [29,30].

4. CONCLUSIONS

Influence of temperature and irradiation by neutrons on the optical absorption spectrum of nickel oxide was studied in the near infra-red energy range 7750-8300 cm⁻¹. This interval corresponds to the magnetic-dipole transition ${}^{3}A_{2g}(F) \rightarrow$ ${}^{3}T_{2g}(F)$ at Ni²⁺ sites. Two sharp lines were observed at the low-energy side of the band: the peak at 7805 cm⁻¹ is assigned to the pure exciton transition, whereas the peak at 7845 cm⁻¹ to the exciton-magnon excitation (magnon sideband), which occurs at the Brillouin zone-centre (BZC). An increase of temperature or of the defects concentration at high neutron fluences results in a lowering of the magnon satellite peak intensity. The effect of temperature leads to a homogenious broadening of the magnon sideband. At the same time, a decrease of the magnon sideband intensity upon irradiation with neutrons is attributed to a distruction of the long-range magnetic ordering due to inhomogenious broadening.

ACKNOWLEDGMENTS

This work was partially supported by grants of the Latvian Government (No. 01.0806 and 01.0821) and the Estonian Science Foundation (No. 3453).

REFERENCES

- 1. S. Hüfner, "Electronic structure of NiO and related 3d-transition-metal compounds," Adv. Phys. 43, pp. 183-356, 1994.
- 2. W.L. Roth, "Magnetic Structures of MnO, FeO, CoO, and NiO," Phys. Rev. 110, pp. 1333-1341, 1958.
- 3. A.Z. Menshikov, Yu.A. Dorofeev, A.G. Klimenko, and N.A. Mironova, "Magnetic phase diagram of (Ni_{1-x}Mg_x)O solid solutions," *Phys. Status Sol. (b)* **164**, pp. 275-283, 1991.
- 4. R.E. Dietz, W.F. Brinkman, A.E. Meixner, and H.J. Guggenheim, "Raman scattering by four magnons in NiO and KNiF₃," *Phys. Rev. Lett.* **27**, pp. 814-817, 1971.

- 5. M.T. Hutchings and E.J. Samuelsen, "Measurement of spin-wave dispersion in NiO by inelastic neutron scattering and its relation to magnetic properties," *Phys. Rev. B* 6, pp. 3447-3461, 1972.
- 6. M.J. Massey, N.H. Chen, J.W. Allen, and R. Merlin, "Pressure dependence of two-magnon Raman scattering in NiO," *Phys. Rev. B* 42, pp. 8776-8779, 1990.
- 7. N.A. Mironova, G.A. Grinvald, V.N. Skvortsova, and U.A. Ulmanis, "Fine structure of absorption spectra of antiferromagnetic NiO," *Sov. Phys. Solid State* 23, pp. 874-875, 1981.
- 8. G.A. Grinvald and N.A. Mironova, "Temperature dependencies of near infrared luminescence from isolated and exchange-coupled pairs of Ni²⁺ ions in MgO," *Phys. Status Sol. (b)* **99**, pp. K101-K102, 1980.
- 9. N.A. Mironova and V.N. Skvortsova, "Fine structure of the NiO:Co infrared absorption spectrum," *Phys. Status Sol. (b)* **127**, pp. K81-K82, 1985.
- N. Mironova, V. Skvortsova, A. Kuzmin, I. Sildos, and N. Zazubovich, "Low temperature optical absorption by magnons in KNiF₃ and NiO single-crystals," in *Defects and Surface-Induced Effects in Advanced Perovskites*, G. Borstel, A. Krumins, and D. Millers, ed., pp. 155-160, Kluwer Academic, Dordrecht, 2000.
- 11. N. Mironova-Ulmane, V. Skvortsova, A. Kuzmin, and I. Sildos, "Exciton-magnon interactions in single-crystal Ni_cMg_{1-c}O solid solutions," *Phys. Solid State* **44**, pp. 1463-1467, 2002.
- 12. A. Kuzmin, N. Mironova-Ulmane and S. Ronchin, "Origin of visible photoluminescence in NiO and Ni_cMg_{1-c}O single-crystals," *Proc. SPIE* **5122**, pp. 61-67, 2003.
- 13. R.E. Dietz, G.I. Parisot, and A.E. Meixner, "Infrared absorption and Raman scattering by two-magnon processes in NiO," *Phys. Rev. B* **4**, pp. 2302-2310, 1971.
- 14. D.J. Lockwood, M.G. Cottam, and J.H. Baskey, "One- and two-magnon excitations in NiO," J. Magn. Magn. Mater. 104–107, pp. 1053-1054, 1992.
- 15. M. Grimsditch, L.E. McNeil, and D.J. Lockwood, "Unexpected behavior of the antiferromagnetic mode of NiO," *Phys. Rev. B* 58, pp. 14462-14 466, 1998.
- 16. Y. Mita, Y. Ishida, M. Kobayashi and S. Endo, "Pressure effects on first-order magnetic Raman scattering in NiO," *J. Phys.: Condens. Matter* 14, 11173–11176, 2002.
- 17. E. Cazzanelli, A. Kuzmin, G. Mariotto, and N. Mironova-Ulmane, "Study of vibrational and magnetic excitations in Ni_cMg_{1-c}O solid solutions by Raman spectroscopy," *J. Phys.: Condens. Matter* **15**, pp. 2045-2052, 2003.
- 18. N.A. Mironova and U.A. Ulmanis, "*Radiation Defects and Metal Ions of Iron Group in Oxides*," Zinatne, Riga, 1988.
- Z. Feng and M.S. Seehra, "Phase diagram and magnetic properties of the diluted fcc system Ni_pMg_{1-p}O," *Phys. Rev. B* 45, pp. 2184-2189, 1992.
- 20. D. Mertz, Y. Ksari, F. Celestini, J. M. Debierre, A. Stepanov, and C. Delmas, "Magnetic cluster formation in Li_xNi_{1-x}O compounds: Experiments and numerical simulations," *Phys. Rev. B* **61**, pp. 1240-1245, 2000.
- 21. A. J. Devasahayama and M. H. Kryder, "The dependence of the antiferromagnet/ferromagnet blocking temperature on antiferromagnet thickness and deposition conditions," J. Appl. Phys. 85, pp. 5519-5521, 1999.
- 22. V. Wagner, D. Tocchetti, and B. Hennion, "Magnetic excitations in antiferromagnetic Ni_{1-x}Co_xO," *AIP Conf. Proc.* **29**, pp. 255-256, 1976.
- 23. F. Haaß, Buhrmester, and M. Martin, "High-temperature in situ X-ray absorption studies on the iron valence in iron-doped nickel oxide (Ni_xFe_{1-x})₁₋₈O," *Solid State Ionics* **141–142**, pp. 289-293, 2001.
- 24. N.A. Mironova and G.V. Bandurkina, "Epitaxial grow of single-crystal Ni_cMg_{1-c}O solid solutions," *Izv. Akad. Nauk Latv. SSR, Ser. Fiz. Tech. Nauk* **4**, pp. 14-19, 1975.
- 25. K.K. Rebane, "Zero-phonon lines in the spectroscopy and photochemistry of impurity-doped solid matter," in *Zero-Phonon Lines and Spectral Hole Burning in Spectroscopy and Photochemistry*, O. Sild and K. Haller, ed., pp.1-18, Springer, Berlin, 1988.
- 26. A.J. Sievers and M. Tinkham, "Far infrared antiferromagnetic resonance in MnO and NiO," *Phys. Rev.* **129**, pp. 1566-1571, 1963.
- 27. P.L. Richards, "Far-infrared magnetic resonance in CoF₂, NiF₂, KNiF₃, and YbIG," *J. Appl. Phys.* **34**, pp. 1237-1238, 1963.

- 28. S.I. Yun, L.A. Kappers, and W.A. Sibley, "Enhancement of impurity-ion absorption due to radiation-produced defects," *Phys. Rev. B* **8**, pp. 773-779, 1973.
- 29. K.W. Blazey, "Optical absorption of MgO:Fe," J. Phys. Chem. Solids 38, pp. 671-675, 1977.
- 30. A. Hjortsberg, J.T. Vallin, and F.S. Ham, "Jahn-Teller effects in the near-infrared absorption spectrum of MgO:Fe²⁺," *Phys. Rev. B* **37**, pp. 3196-3215, 1988.